

The Principle of Complete Positivity and the Open Systems View

(based on joint work with Stephan Hartmann)

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Complete positivity:

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Our claim:

- Complete positivity is an expression of what we call the “closed systems view of quantum theory.”
- We should reject the closed systems view (in favour of the “open systems view”) and deny complete positivity the status of a fundamental physical principle.

Outline:

1. The principle of complete positivity
2. The closed systems view vs. the open systems view
3. Making sense of fundamental open systems dynamics in Everettian and informational interpretations of QM
4. Conclusion

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- Probabilistic generalisation of a state vector:

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- Compound system (product state):
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- State change of \mathcal{S} in the presence of \mathcal{E} :

$$\rho_{\mathcal{S}} \mapsto \text{tr}_{\mathcal{E}}(\mathbf{U}\rho_{\mathcal{S}} \otimes \rho_{\mathcal{E}}\mathbf{U}^\dagger) = \rho'_{\mathcal{S}} = \Lambda\rho_{\mathcal{S}}$$

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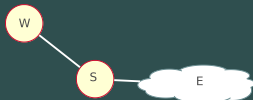
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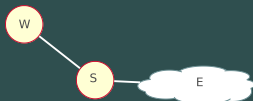
Consider the system $\mathcal{S} + \mathcal{W}_n$ evolving in the presence of \mathcal{E} , such that:

- \mathcal{W}_n : A system of dimensionality n not currently interacting with \mathcal{S} , but which may have interacted with it in the past.
- \mathcal{W}_n evolves trivially.

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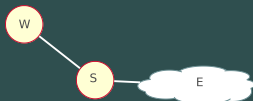
Problem: Requiring that Λ be positive on \mathcal{S} does not guarantee that $\Lambda \otimes I_n$ is positive on $\mathcal{S} + \mathcal{W}_n$.

- I.e., $\Lambda \otimes I_n$ will map some of the states in $\mathcal{H}_{\mathcal{S}+\mathcal{W}_n}$ to unphysical states (i.e., that yield negative probabilities for the results of measurements on $\mathcal{S} + \mathcal{W}_n$).

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Solution: Require that $\Lambda \otimes I_n$ be positive for all \mathcal{W}_n .

“One may reasonably doubt this argument. It is very powerful magic: \mathcal{W} sits apart from $\mathcal{S} + \mathcal{E}$ and does absolutely nothing; by doing so, it forces the motion of \mathcal{S} to be completely positive with dramatic physical consequences ...” (Pechukas, 1994).

More concrete reasons to be skeptical (Shaji & Sudarshan, 2005):

Suppose, on the one hand, that \mathcal{S} and \mathcal{W}_n are not entangled:

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- If we enforce complete positivity then it would seem to follow that no valid physical description of the dynamics of \mathcal{S} can be given when it is initially entangled with \mathcal{E} .

Not completely positive (NCP) maps:

It is precisely for setups like these that an NCP map will make sense.

- When \mathcal{S} and \mathcal{E} are entangled, then it follows that it is impossible for \mathcal{S} to be (for instance) in a pure state, or in general in any state that is not a valid partial trace over the combined state of $\mathcal{S} + \mathcal{E}$.

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- Such states that are outside of the 'compatibility domain' of an NCP map will be ill-described by it.
- But as long as such a map is completely positive **in relation to all of the actually possible states of \mathcal{S}** in a given setup, it seems that there is no reason not to use it to describe the dynamics of \mathcal{S} (Cuffaro & Myrvold, 2013, sec. 5).

A better argument for imposing complete positivity as a fundamental physical principle:

A system-theoretic description of an open system has to be considered as phenomenological; **the requirement that it should be derivable from the fundamental automorphic dynamics of a closed system** implies that the dynamical map of an open system has to be completely positive. (Raggio & Primas, 1982, p. 435, our emphasis).

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Standard quantum theory (ST)

- Physical state of a closed system, \mathcal{S} , is represented by a state vector, $|\psi\rangle$, in a Hilbert space for \mathcal{S} .
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Strictly speaking no system (except, perhaps, the whole universe) can really be isolated. How do we model open systems in **ST**?

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Closed systems view of an open system:

- \mathcal{S} 's interaction with its environment is described in terms of its being coupled with a separate system \mathcal{E} such that $\mathcal{S} + \mathcal{E}$ form an isolated system.



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- Guaranteed by Stinespring's dilation theorem (Stinespring, 1955), which assumes complete positivity.

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- Black hole physics gives us (prima facie) reasons to motivate describing the evolution of the cosmos as formally similar to the evolution of an open system (Hawking, 1976).
- Global unitary evolution is hard to square with important recent approaches to quantum gravity (Oriti, 2021, sec. 3.1).

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- **It isn't the dynamics of $\mathcal{S} + \mathcal{E}$, but the dynamics of \mathcal{S} , that we take ourselves to have successfully described when we do this.**¹
- Thus there is a clear **empirical motivation** to extrapolate from the dynamics of open systems rather than from the dynamics of closed systems.

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Cuffaro & Hartmann:

The Open Systems View (arXiv:2112.11095)

- Open systems, represented by density operators, are fundamental in quantum theory.



Cf. Chen (2018).

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- Λ_t acts on the state space of \mathcal{S} (not on $\mathcal{S} + \mathcal{E}$).
- Formulated in accordance with the open systems view
 - The environment is not represented as a separate system; its influence is represented in the dynamical equations that govern the evolution of \mathcal{S} .

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So GT, unlike ST, allows for **fundamental non-unitary evolution**, and is, in this sense, a more general dynamical framework than ST (despite not really adding anything to the Hilbert space formalism).

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 - Both take (vanilla) quantum mechanics as providing the resources with which to completely describe reality.
 - But in terms of their a priori commitments, these two interpretations couldn't be more different.
 - Yet, we argue, both have very good reasons to endorse the open systems view of quantum theory.

$$\begin{aligned} |\psi\rangle_S &= \alpha|b_1^+\rangle + \beta|b_1^-\rangle \\ &= \alpha'|b_2^+\rangle + \beta'|b_2^-\rangle. \end{aligned}$$

What does this mean on the informational (neo-Bohrian) interpretation?

- Coupling the degrees of freedom of \mathcal{S} to those of a further system \mathcal{M} will yield a collection of unitarily-related conditional probability distributions over the possible outcomes of an assessment of \mathcal{M} as described with respect to a particular basis b_m .

“In the treatment of atomic problems, actual calculations are most conveniently carried out with the help of a Schrödinger state function, from which the statistical laws governing observations obtainable under specified conditions can be deduced by definite mathematical operations. It must be recognized, however, that we are here dealing with a purely symbolic procedure, **the unambiguous physical interpretation of which in the last resort requires a reference to a complete experimental arrangement.**” (Bohr, 1958, pp. 392–393, our emphasis).

Informational (neo-Bohrian) interpretation:

- Notice that \mathcal{S} is conceived of here as an open system (even when its state is described by a state vector), **but since open systems dynamics are not fundamental in ST, we require a larger Hilbert space** (including the degrees of freedom of both \mathcal{S} and \mathcal{M}) to represent it as such.

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- **ST is about open systems**, on the informational interpretation, despite being formulated from the closed systems view (which it inherits from classical mechanics).
- The motivation for adopting GT, which takes open systems to be fundamental, as our theoretical framework for describing reality is very clear.

What about the universe as a whole?

- Note that it isn't essential that a particular system represented by \mathcal{M} actually exists.
- What's important is the dynamical context that \mathcal{M} represents. It's always possible to imagine a dynamical physical interaction with the empirically accessible degrees of freedom of any physical system—because, conceptually, that is just what we mean when we say that a physical system is empirically accessible—regardless of that system's size (Janas, Cuffaro, & Janssen, 2022, p. 216).
- The further question of whether it “really” makes sense to talk about the empirical accessibility of the universe as a whole is a philosophical question. But as far as physics is concerned, the idea of empirical accessibility is given a clear meaning, in the framework of QM, that can in principle be applied to *any* physical system.

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- Fundamental unitarity?
 - Not necessarily! (see, e.g., Wallace 2012, sec. 10.5).
 - Ultimately the Everett interpretation is committed to **quantum theory**.
 - But GT **makes no changes** to the quantum formalism. Rather, it takes a different view of that formalism—one that allows for fundamental non-unitary evolution.

Many worlds:

$$\rho = p |\psi_1\rangle\langle\psi_1| + (1 - p) |\psi_2\rangle\langle\psi_2|.$$

- The fact that the different terms, $|\psi_1\rangle\langle\psi_1|$ and $|\psi_2\rangle\langle\psi_2|$ are by definition **decoherent** makes it unproblematic, irrespective of whether ρ evolves unitarily, to identify them with independently evolving worlds.
- This is even clearer than the FAPP story one needs to give in the pure state, unitary, case.

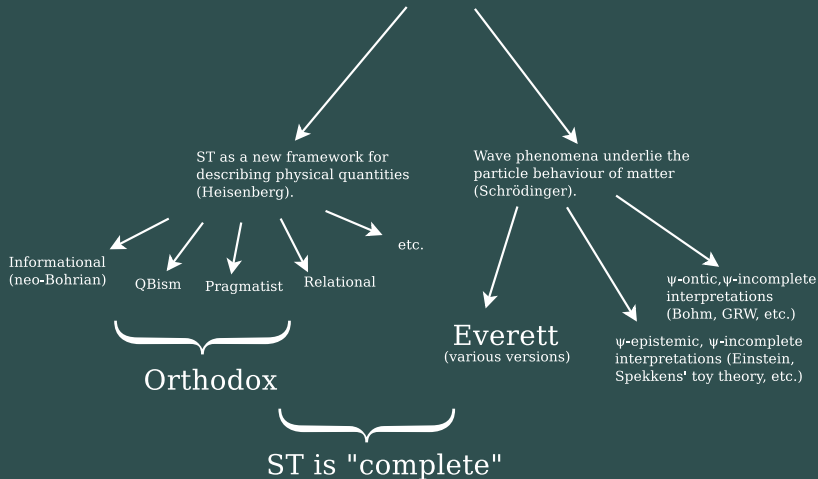
Upshot:

- Fundamental open systems dynamics can be physically motivated.
- There is as much empirical support, if not more, for GT as there is for ST.
- GT is as much in the spirit of Everett as ST is.

There seems to be little reason for an Everettian not to embrace GT as the proper dynamical framework with which to engage in foundational and philosophical investigations (and speculations) regarding the quantum world.

Interpretations of ST

(Janas et al, 2022)



Outline:

1. The principle of complete positivity
2. The closed systems view vs. the open systems view
3. Making sense of fundamental open systems dynamics in Everettian and informational interpretations of QM
4. Conclusion

~~A system-theoretic description of an open system has to be considered as phenomenological; the requirement that it should be derivable from the fundamental automorphic dynamics of a closed system implies that the dynamical map of an open system has to be completely positive. (Raggio & Primas, 1982, p. 435).~~

Conclusion: We should reject complete positivity as a fundamental physical principle.

Works Cited I

- Bohr, N. (1958). Quantum physics and philosophy. In R. Klibansky (Ed.) *Philosophy in the Mid-Century: A Survey*, (pp. 308–314). Firenze: La Nuova Italia Editrice. Page reference to reprint in: J. Kalckar, Ed., *Niels Bohr: Collected Works*, Vol. 7 (pp. 385–394). Amsterdam: Elsevier.
- Chen, E. K. (2018). Quantum mechanics in a time-asymmetric universe: On the nature of the initial quantum state. *The British Journal for the Philosophy of Science*. In press.
- Cuffaro, M. E., & Hartmann, S. (2021). The open systems view. arXiv:2112.11095v1.
- Cuffaro, M. E., & Myrvold, W. C. (2013). On the debate concerning the proper characterisation of quantum dynamical evolution. *Philosophy of Science*, 80, 1125–1136.
- Curiel, E. (2019). On geometric objects, the non-existence of a gravitational stress-energy tensor, and the uniqueness of the Einstein field equation. *Studies in History and Philosophy of Modern Physics*, 66, 90–102.
- Gryb, S., & Sloan, D. (2021). When scale is surplus. arXiv:2103.07384v2.
- Hawking, S. W. (1976). Breakdown of predictability in gravitational collapse. *Physical Review D*, 14, 2460–2473.
- Hofer, C. (2000). Energy conservation in GTR. *Studies in History and Philosophy of Modern Physics*, 32(2), 187–199.

Works Cited II

- Janas, M., Cuffaro, M. E., & Janssen, M. (2022). *Understanding Quantum Raffles: Quantum Mechanics on an Informational Approach: Structure and Interpretation*. Springer-Verlag. In press.
- Jordan, T. F., Shaji, A., & Sudarshan, E. C. G. (2004). Dynamics of initially entangled open quantum systems. *Physical Review A*, *70*, 052110.
- Maudlin, T., Okon, E., & Sudarsky, D. (2020). On the status of conservation laws in physics: Implications for semiclassical gravity. *Studies in History and Philosophy of Modern Physics*, *69(1)*, 67–81.
- Oriti, D. (2021). The complex timeless emergence of time in quantum gravity. In P. Harris, & R. Lestienne (Eds.) *Time and Science*. World Scientific. Forthcoming.
- Pechukas, P. (1994). Reduced dynamics need not be completely positive. *Physical Review Letters*, *73*, 1060–1062.
- Raggio, G. A., & Primas, H. (1982). Remarks on “On completely positive maps in generalized quantum dynamics”. *Foundations of Physics*, *12*, 433–435.
- Shaji, A., & Sudarshan, E. C. G. (2005). Who's afraid of not completely positive maps? *Physics Letters A*, *341*, 48–54.
- Sloan, D. (2018). Dynamical similarity. *Physical Review D*, *97*, 123541.
- Smeenk, C., & Ellis, G. (2017). Philosophy of cosmology. In E. N. Zalta (Ed.) *The Stanford Encyclopedia of Philosophy*. Metaphysics Research Lab, Stanford University, Winter 2017 ed.
- Stinespring, W. F. (1955). Positive functions on C*-algebras. *Proceedings of the American Mathematical Society*, *6*, 211.
- Wallace, D. (2012). *The Emergent Multiverse*. Oxford: Oxford University Press.